Lecture 10 : Rigid Body III.

- Rotation about a moving axis.
- Rolling without slipping.
- Rolling friction
- Work & Power

So far we know:

① Translation motion of CM : $\sum \vec{F}_{ext} = m \vec{a}_{cm}$

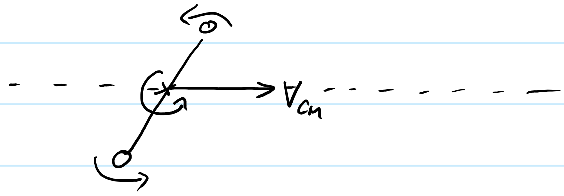
② Rotation of rigid body about a fixed axis : $\sum \tau_{ext_z} = I_z \alpha_z$

Combine both types of motion : a rigid body rotating about a moving axis.

e.g. Rolling wheel.

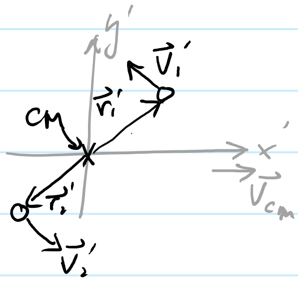
$\boxed{\text{Translation motion of CM + Rotational motion about a axis through its CM.}}$

Energy.



In CM-frame

Pure rotation



choose

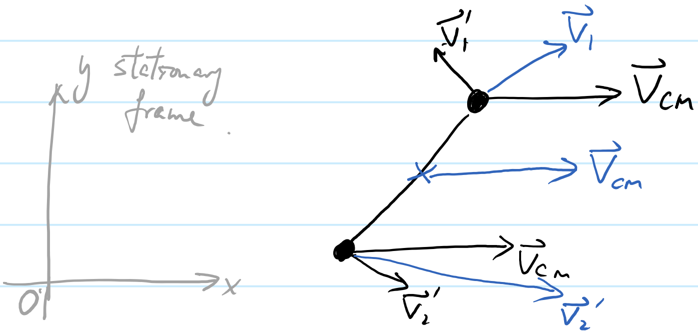
$$\vec{0} = \vec{r}_{cm}' = \frac{m_1 \vec{r}_1' + m_2 \vec{r}_2'}{m_{tot}}$$

$$\frac{d}{dt} : \vec{0} = \vec{V}_{cm}' = \frac{m_1 \vec{v}_1' + m_2 \vec{v}_2'}{m_{tot}}.$$

velocity of cm in CM frame

Add $\vec{V}_{cm}$ to all points on the object.

i.e. add a translational motion.



$\vec{V}_i$: velocity measured by a stationary reference frame

In a stationary frame,

$$K = \sum_i \frac{1}{2} m_i |\vec{V}_i|^2$$

$$= \frac{1}{2} \sum m_i \left(\vec{V}_i \cdot \vec{V}_i\right)$$

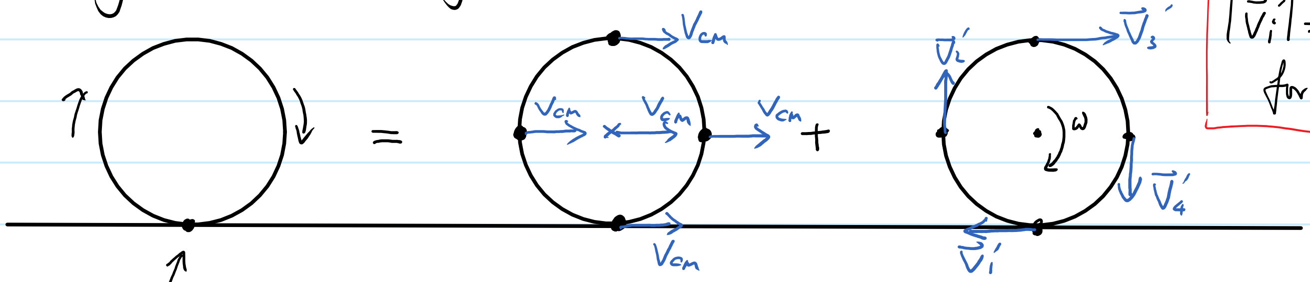
$$= \frac{1}{2} \sum m_i \left(\vec{V}_i' + \vec{V}_{cm}\right) \cdot \left(\vec{V}_i' + \vec{V}_{cm}\right)$$

$$= \frac{1}{2} \sum m_i \left|\vec{V}_i'\right|^2 + \frac{1}{2} \sum m_i \left|\vec{V}_{cm}\right|^2 + \frac{1}{2} \sum m_i \left(2 \vec{V}_{cm} \cdot \vec{V}_i'\right)$$

$$= \frac{1}{2} \sum m_i (r_i \omega)^2 + \frac{1}{2} m_{tot} V_{cm}^2 + \vec{V}_{cm} \cdot \underbrace{\sum m_i \vec{V}_i'}_{\vec{V}_{cm}' = \vec{0}}$$

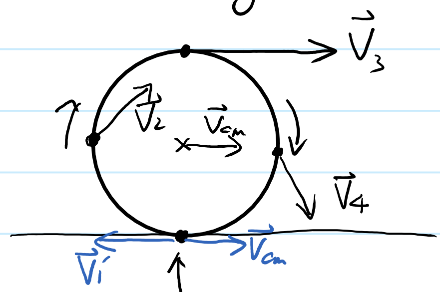
$$\boxed{K = \underbrace{\frac{1}{2} m_{tot} V_{cm}^2}_{K_{tran}} + \underbrace{\frac{1}{2} I \omega^2}_{K_{rot,cm}}}$$

Rolling without slipping



$|\vec{V}_i'| = R\omega$ for all i.

stationary   translation   +   rotation about CM.



$\vec{V}_1 = \vec{V}_{cm} + \vec{V}_1'$

Velocity of each point relative to the ground

$$\vec{V}_i = \vec{V}_i' + \vec{V}_{cm}$$

Velocity relative to CM + Velocity of CM

Require: $\vec{V}_1 = \vec{0}$ stationary point.

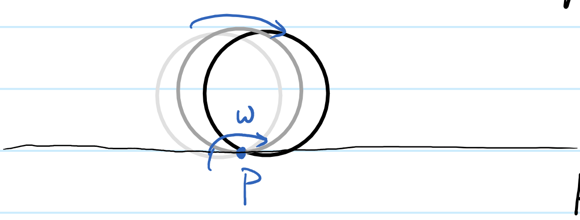
$\Rightarrow \quad \vec{V}_1' = -\vec{V}_{cm}$

$$\underline{V_{cm} = V_1' = R\omega}$$

Kinematic constrain for rolling without slipping.

Alternative view on rolling.

Rolling = swinging about the point of contact "P" but the contact point changes continuously.



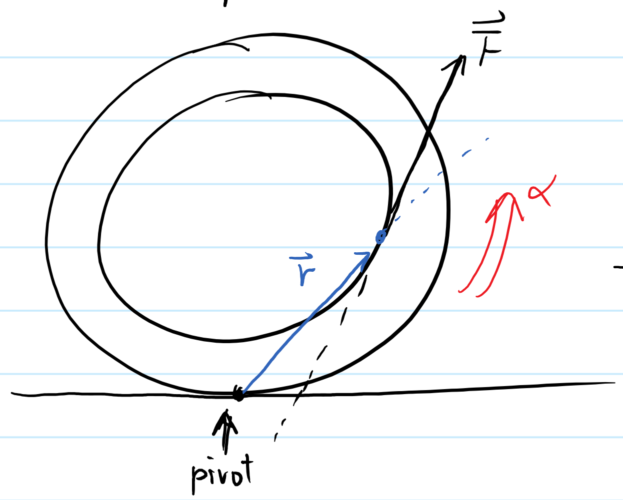
$K = \frac{1}{2} I_p \omega^2$, $I_p = I_{cm} + MR^2$

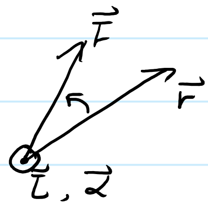
$K = \frac{1}{2} I_{cm} \omega^2 + \frac{1}{2} MR^2 \cdot \omega^2$

$= \frac{1}{2} I_{cm} \omega^2 + \frac{1}{2} M V_{cm}^2 \quad \because R\omega = V_{cm}$

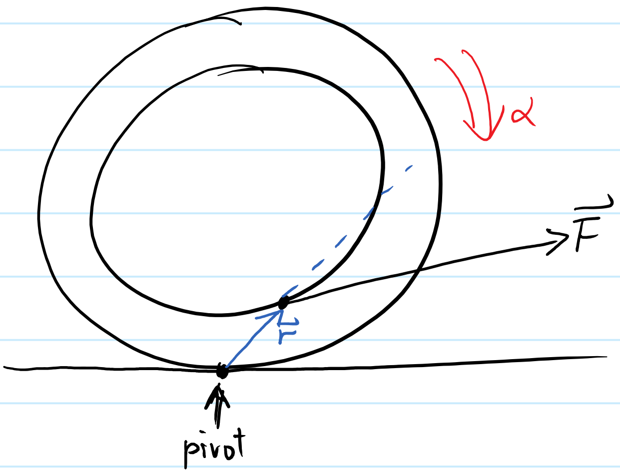
same as previous page

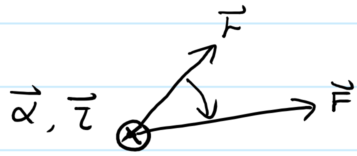
Demo. spool.





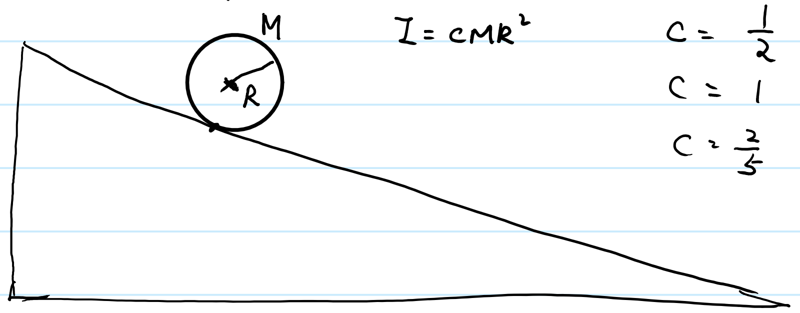
$\vec{r} \times \vec{F}$: out of page $\Rightarrow$ $\vec{\alpha}$ out of page





$\vec{r} \times \vec{F}$: into the page $\Rightarrow$ $\vec{\alpha}$ into the page

Rolling down inclined plane.



$I = cMR^2$

$c = \frac{1}{2}$ cylinder
$c = 1$ ring
$c = \frac{2}{5}$ sphere

contact point does not move $\Rightarrow$ static friction $\Rightarrow W_{friction} = 0$

$$E_i = E_f$$

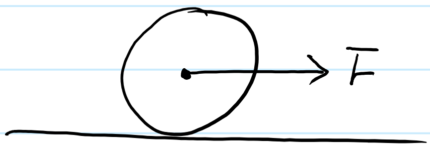
$$Mgh = \frac{1}{2} M V_{cm}^2 + \frac{1}{2} I_{cm} \omega^2 \qquad R\omega = V_{cm}$$

$$Mgh = \frac{1}{2} M V_{cm}^2 + \frac{1}{2} \frac{cMR^2}{R^2} V_{cm}^2$$

$$V_{cm} = \sqrt{\frac{2gh}{1+c}}$$

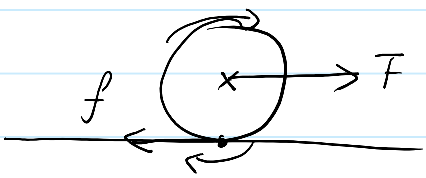
depends only on type of rolling object, c.
not M or R.

Torque Approach.

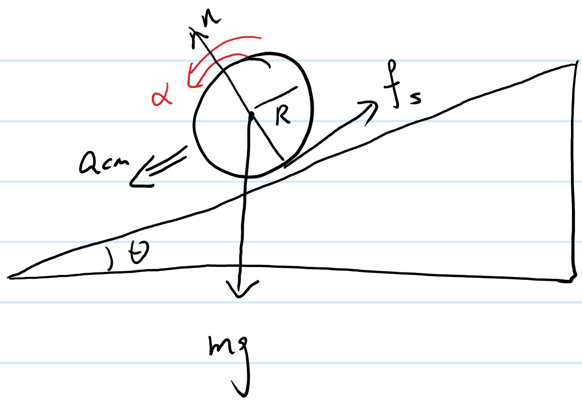


no friction $\Rightarrow$ no torque about cm
$\Rightarrow$ no angular acceleration
$\Rightarrow$ only sliding occurs.

To make it rolls,



If it is not slipping $\Rightarrow$ contact point not sliding
$\Rightarrow$ static friction.
f_s.



A cylinder rolling down without slipping

given m, R.,
solve for a_{cm} & α.

Translation motion of CM governed by Newton's 2nd Law.

$$\sum \vec{F}_{ext} = m\vec{a}_{cm}.$$

Rotational motion about CM governed by

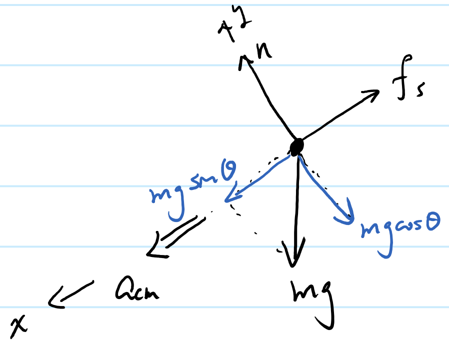
$$\sum \tau_{ext_{cm}} = I_{cm}\,\alpha$$

In general, $\vec{a}_{cm}$ and α are independent of each other.

But for rolling without slipping, $\underline{V_{cm} = R\omega} \Rightarrow \underline{a_{cm} = R\alpha}$.

We have

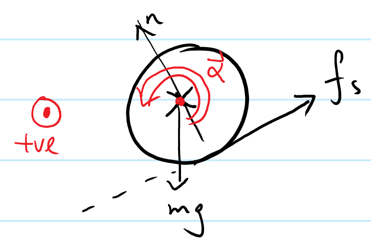
<u>For translation motion</u>



$$y: \quad n - mg\cos\theta = 0 \quad \text{——} \ ①$$

$$x: \quad mg\sin\theta - f_s = m a_{cm} \quad \text{——} \ ②$$

<u>For rotational motion about CM</u>



only f_s produce a torque about cm.

$$\tau_{cm} = I_{cm}\alpha$$

$$+f_s \cdot R = +I_{cm}\,\alpha \quad \text{——} \ ③$$

Kinematic constrain: $a_{cm} = R\alpha$ — ④

$$\Rightarrow \begin{cases} n = mg\cos\theta & \text{— ①} \\ mg\sin\theta - f_s = m a_{cm} & \text{— ②} \\ f_s R = I_{cm}\alpha & \text{— ③} \\ a_{cm} = R\alpha & \text{— ④} \end{cases}$$

$\Rightarrow$ eliminating f_s and α, ② ③ & ④ gives.

$$\Rightarrow mg\sin\theta - I_{cm}\frac{a_{cm}}{R^2} = m a_{cm}$$

$$a_{cm} = \frac{g\sin\theta}{1 + \frac{I_{cm}}{mR^2}}$$

$$\Rightarrow a_{cm} = \frac{2}{3} g\sin\theta \qquad \because I_{cm} = \frac{1}{2}mR^2 \text{ for cylinder.}$$

$$\alpha = \frac{a_{cm}}{R} = \frac{2}{3}\frac{g\sin\theta}{R}$$

Summary. To solve dynamical problem with both translational and rotational motions, we can use.

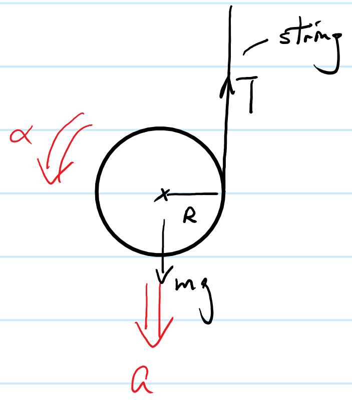
$\begin{cases} \text{Newton's 2nd Law for the trans. motion of its cm.} \\ + \\ \text{Newton's 2nd Law for rotation for the rotational motion about its cm.} \end{cases}$

i.e. $\boxed{\Sigma \vec{F}_{ext} = m a_{cm}}$ & $\boxed{\Sigma \tau_{cm} = I_{cm}\alpha}$

For $\boxed{\text{Rolling without slipping,}}$ we add the

kinematic constrain: $V_{cm} = R\omega$, $a_{cm} = R\alpha$

Yo-yo : vertical rolling without slipping along a string.



assume the yo-yo is a solid cylinder for simplicity.

$$I_{cm} = \frac{1}{2} m R^2$$

Translational motion: $\Sigma F = mg - T = ma$ ——— ①

Rotational motion about CM: $\Sigma \tau_{cm} = T \cdot R = I_{cm} \alpha$ ——— ②

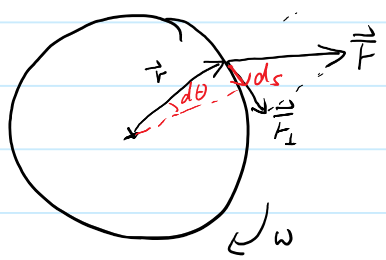
Kinematic constrain: $\alpha = a/R$ ——— ③

$$\Rightarrow \quad mg - \frac{I_{cm}\alpha}{R} = ma$$

$$\Rightarrow \quad mg - \frac{I_{cm} a}{R^2} = ma$$

$$\Rightarrow \quad \begin{cases} a = \frac{2}{3} g \\ T = \frac{1}{2} ma = \frac{1}{3} mg \end{cases}$$

Work and Power



$$dW = \vec{F} \cdot d\vec{s}$$
$$= F_\perp \cdot ds$$
$$= F_\perp \cdot r\, d\theta$$
$$= \tau\, d\theta$$
$$\Rightarrow \boxed{W = \int \tau\, d\theta}$$

$$\boxed{P = \vec{\tau} \cdot \vec{\omega}}$$

c.f. $W = \int \vec{F} \cdot d\vec{s}$ & $P = \vec{F} \cdot \vec{v}$

Rolling without slipping

- contact point does not slide.
- friction on the wheel is static friction.

$\Rightarrow$ $W_{friction} = 0$ static friction does not do work.

$\Rightarrow$ No energy loss!

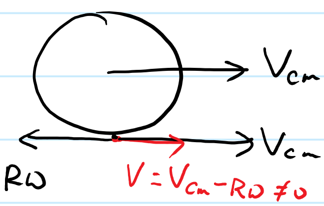
$\Rightarrow$ Roll forever if no other resistances.

But, in reality, rolling without slipping is rare.

The surface of the road is not perfectly flat.

If the wheel hits a bump, the wheel could jump up. When it touches the ground again, its V_{cm} and ω may not satisfy $V_{cm} = R\omega$ anymore.

If $V_{cm} \neq R\omega$, the point of the wheel contacting the ground will slide.



The sliding motion will introduce a kinetic friction which leads to energy loss.